A Note on Gauge Transformations and Stabile Ordinary Differential Equations

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1 Introduction

We have introduced the notion of stabile ODEs (odrinary differential equations) for obtaining numerical solutions valid when the independent variable x goes to the infinity. In our applications of holonomic gradient method (HGM, see, e.g., [2], [3], [4]), it is a key step to find a stabile ODE [1]. In this short note, we will propose an algorithmic method, which utilizes Gauge transformations, to derive a stabile system. If we can apply this algorithmic method by an algebraic calculation, we can obtain a stabile system with no numerical error. We demonstrate that this approach gives a useful stabile ODE used in [1].

2 Obtaining Stabile Subsystems

We review our definition of a stabile ODE given in [1].

Definition 1. Consider the ODE

$$\left[a_m(x)\frac{d^m}{dx^m} + \dots + a_1(x)\frac{d}{dx} + a_0(x)\right] \bullet u = 0$$
(1)

where $a_i(x)$ are polynomials in x. We assume that f is a solution of this differential equation. and denote with $f_1(x), \ldots, f_m(x)$ its linearly independent solutions. Then, let $f_i(x)$ be the dominant solution for $x \to \infty$, i.e., $|f_i(x)| \ge |f_j(x)|$, $\forall j$. We refer to the LODE (liear ODE) as *stabile* for f(x) if $\lim_{x\to\infty} \frac{|f_i(x)|}{|f(x)|} < \infty$. Note that a LODE is stabile or not regardless of the selected set of linearly independent solutions. The notion of stabile LODE is defined analogously in the case of a vector-valued function, by replacing $|\cdot|$ with a vector norm $||\cdot||$.

Theorem 1. Stabile LODEs can be derived algorithmically and numerically as a lower-dimensional subsystem of given LODEs.

Proof. We give our algorithm. It is enough to give an algorithm in case of a vector valued ODE (a system of ODEs) of the first order.

Let $\mathbf{g} = (g_1, \dots, g_r)^{\mathsf{T}}$ be the dominant solution of the LODE $\mathbf{f}' = \mathbf{P}(x)\mathbf{f}$ where \mathbf{P} is an $r \times r$ square matrix valued function. For this LODE, let us consider the following gauge transformation matrix and its inverse:

$$\mathbf{G} = \begin{pmatrix} g_1 & 0 & 0 & \cdots & 0 \\ g_2 & & \\ \vdots & & \\ g_r & & \\ g_r & & \\ \end{pmatrix}, \quad \mathbf{G}^{-1} = \begin{pmatrix} g_1^{-1} & 0 & 0 & \cdots & 0 \\ g_2^{-1} & & \\ \vdots & & \\ \vdots & & \\ -g_r g_1^{-1} & & \\ \mathbf{I}_r - 1 \end{pmatrix}. \quad (2)$$

Then, the matrix of the transformed LODE by the Gauge transformation ${\bf G}$ can be written as

$$\mathbf{G}^{-1}\mathbf{P}\mathbf{G} - \mathbf{G}^{-1}\mathbf{G}' = \begin{pmatrix} 0 & g_1^{-1}p_{12} & g_1^{-1}p_{13} & \cdots & g_1^{-1}p_{1r} \\ 0 & & & \\ \vdots & & & \mathbf{P}_2 \\ 0 & & & & \end{pmatrix},$$
(3)

where $p_{ij} = [\mathbf{P}]_{i,j}$, and \mathbf{P}_2 is an $(r-1) \times (r-1)$ matrix with column (i-1) given by

$$(-g_2g_1^{-1}p_{1i} + p_{2i} - g_3g_1^{-1}p_{1i} + p_{3i} \dots - g_rg_1^{-1}p_{1i} + p_{ri})^{\mathsf{T}}.$$
 (4)

In other words, the original ${\bf P}$ is block upper-triangularized by the gauge transformation ${\bf G}.$

It follows from the local theory of LODEs that the second dominant solution of the LODE for \mathbf{F} is obtained from the dominant solution of the smaller system $\mathbf{f}'_2 = \mathbf{P}_2 \mathbf{f}_2$. We can apply this procedure recursively. //

We apply this procedure for the following ODE discussed in [1]

$$\frac{\partial}{\partial x}\mathbf{f}(x,\lambda) = \mathbf{P}(x,\lambda)\mathbf{f}(x,\lambda) \quad \text{where } \mathbf{f} \text{ is the 4-dimensional vector}$$
(5)

and

$$\mathbf{P}(x,\lambda) = \frac{1}{x\lambda} \begin{pmatrix} a_1 & a_2 & -1 & 0\\ 0 & a_3 & a_2 + 1 & -1\\ a_5 & a_6 & a_7 & n-1\\ a_8 & a_9 & a_{10} & a_{11} \end{pmatrix}$$

 a_i 's are polynomials in x and λ shown in Table 1.

In this case, we can derive a stabile subsystem

$$\frac{\partial}{\partial \varphi} \mathbf{g}(\varphi) = \frac{1}{\varphi} \begin{pmatrix} 0 & 2e^{-\varphi^2 + 2\varphi\psi}\varphi^{2(k+1)} & 0\\ 0 & 0 & 1\\ 0 & -2(2n-1)\varphi\psi & -[4\varphi\psi + 2(n-1)] \end{pmatrix} \mathbf{g}(\varphi).$$
(6)

by an algebraic calculation.

Theorem 2. A gauge transformation yields the 3 dimensional subsystem (2), which is stabile, from the 4 dimensional system (5), which is not stabile, by block upper-triangularization where $x = \sqrt{x}$ and $\psi = \sqrt{\lambda}$.

Table 1: Elements of matrices $\mathbf{P}(x,\lambda)$ from [1] Expression $(k+1)\lambda$ a_1 $\lambda - n + 1$ a_2 $(k+1)\lambda + n - 1$ a_3 $(k+1)x\lambda^2$ a_5 $x\lambda^2 - (n-1)\left[x\lambda + (k+1)\lambda + n - 1\right]$ a_6 $-(x+n-1)\lambda + (n-1)(n-2)$ a_7 $-(n-2)(k+1)x\lambda^2$ a_8 $\begin{array}{l} (k-n+3)x\lambda^2 + (n-1)^2 \left[x\lambda + (k+1)\lambda + n - 1 \right] \\ x\lambda^2 + (n-1)^2 (\lambda - n + 2) \\ -x\lambda - (n-1)^2 \end{array}$ a_9 a_{10} a_{11}

Proof. Our proof relies on [1, Appendix I-b] and we use symbols l_i defined in this appendix. Let \mathcal{I} be the left ideal in ring \mathcal{R} generated by operators l_1 and l_2 respectively, and \mathcal{G} be the Gröbner basis of \mathcal{I} for the lexicographic ordering $\partial_{\lambda} > \partial_x$, which comprises differential operators l_3 , l_4 , and l_5 . Based on the ordering we can identify the leading terms of l_3 , l_4 , and l_5 as θ_x^3 , $\theta_\lambda \theta_x$, and θ_λ^2 , respectively, which, in the (θ_x, θ_λ)-plane, correspond to points (3,0), (1,1), and (0,2), respectively. Then, the remaining points in the top-right quadrant under the line connecting these points are: (0,0), (1,0), (2,0), and (0,1). These points yield the so-called *standard monomials* with respect to \mathcal{G} as 1, θ_x , θ_x^2 , and θ_λ }, respectively. Please see [5], [3, p. 46] for their definition. Finally, notice again that operator l_3 is free of θ_λ .

Then, using the definition of $\theta_x = x \frac{\partial}{\partial x}$ and expressing θ_x^3 from l_3 , we can write the following in the ring of differential operators l_1 and l_2^{-1} :

$$\frac{\partial}{\partial x}1 = \frac{1}{x}\theta_x,\tag{7}$$

$$\frac{\partial}{\partial x}\theta_x = \frac{1}{x}\theta_x^2, \tag{8}$$

$$\frac{\partial}{\partial x}\theta_x^2 = \frac{1}{x}\theta_x^3 \equiv -\frac{2x-2k-3}{x}\theta_x^2 - \frac{(x-k-1)(x-k+n-2)-x\lambda}{x}\theta_x.(9)$$

Because the right-hand side contains θ_x , θ_x^2 , but not θ_λ , we can recast the ODEs in block-diagonalized form as

$$\frac{\partial}{\partial x} \begin{pmatrix} \theta_{\lambda} \\ 1 \\ \theta_{x} \\ \theta_{x}^{2} \end{pmatrix} \equiv \begin{pmatrix} \ast & \ast & \ast & \ast \\ 0 & \mathbf{P}_{2} \end{pmatrix} \begin{pmatrix} \theta_{\lambda} \\ 1 \\ \theta_{x} \\ \theta_{x}^{2} \end{pmatrix} \quad \text{modulo}\,\mathcal{I}, \tag{10}$$

where symbols * stand for unimportant elements, and \mathbf{P}_2 is a 3×3 matrix whose elements can be readily identified from (7), (8), (9).

Now, let \mathcal{J} be the left ideal generated by \mathcal{I} in the ring of differential operators $\mathcal{K}\langle\partial_x,\partial_\lambda\rangle$, where \mathcal{K} is the quotient field of holomorphic functions of (x,λ) on a

 $^{{}^{1}}A \equiv B$ or $A \equiv B$ modulo \mathcal{I} means that $A - B \in \mathcal{I}$. Note that if l_1 and l_2 annihilate f then $A \bullet f = B \bullet f$ when $A \equiv B$.

neighborhood of (1, 1). Using the following gauge transformation²

$$s_1(x,\lambda) = \exp(-2\sqrt{x\lambda})x^{-k-1}e^x\theta_x, \qquad (11)$$

$$s_2(x,\lambda) = \frac{1}{2}\theta_x(\exp(-2\sqrt{x\lambda})x^{-k-1}e^x\theta_x), \qquad (12)$$

we can transform (10) into

$$\frac{\partial}{\partial x} \begin{pmatrix} \theta_{\lambda} \\ 1 \\ s_{1} \\ s_{2} \end{pmatrix} \equiv \begin{pmatrix} \frac{\ast | \ast \ast \ast \ast}{0} \\ 0 \\ 0 \\ 0 \\ \end{bmatrix} \begin{pmatrix} \theta_{\lambda} \\ 1 \\ s_{1} \\ s_{2} \end{pmatrix} \mod \mathcal{J}, \quad (13)$$

where \mathbf{P}_3 is a 3×3 matrix.

Further, let \mathcal{G}' be a Gröbner basis of \mathcal{I} with the ordering $\partial_x > \partial_\lambda$. Recall that this basis is characterized by the differential operators l_6 and l_7 shown in [1]. Computing the *normal forms*[5][3, p. 283] of s_1 and s_2 by \mathcal{G}' yields $c_i(x, \lambda)$ and $d_i(x, \lambda)$, $i = 0, \ldots, 3$, for

$$s_1 \equiv \sum_{i=0}^{3} c_i(x,\lambda) \theta_{\lambda}^i \mod \mathcal{J},$$
 (14)

$$s_2 \equiv \sum_{i=0}^{3} d_i(x,\lambda) \theta^i_{\lambda} \mod \mathcal{J}.$$
 (15)

Now, using $c_i(x, \lambda)$ and $d_i(x, \lambda)$, i = 0:3 to construct the matrix

$$\mathbf{G} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ c_0 & c_1 & c_2 & c_3 \\ d_0 & d_1 & d_2 & d_3 \end{pmatrix},$$
(16)

and applying the gauge transformation with **G** to (5) yields the matrix \mathbf{P}_3 in (13). Finally, upon making the variable transformation $\varphi = \sqrt{x}$, differential equation $\frac{\partial}{\partial x} \mathbf{f} = \mathbf{P}_3 \mathbf{f}$ yields (6). //

Numerical evidences that the stabile ODE (2) is useful in a performance evaluation of multiple-input multiple-output (MIMO) wireless communications systems under Rician fading channel is given in [1].

References

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²Conversely, operators θ_x and θ_x^2 can readily be expressed in terms of s_1 and s_2 .

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