# ELLIPTIC RUIJSENAARS OPERATORS AND ELLIPTIC HYPERGEOMETRIC INTEGRALS 

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#### Abstract

We study a family of mutually commutative difference operators associated with root systems and discuss their simultaneous eigenvectors in a special case. For root systems with rank $n$, we construct $3 n$ commutative difference operators, which are a generalization of elliptic Ruijsenaars operators. In particular, for the $B C_{1}$ root system, we construct an explicit simultaneous eigenvector of these operators described in terms of elliptic hypergeometric integrals.


## 1. Introduction

In [14] Ruijsenaars introduced the operators acting on the space of meromorphic functions which are defined by

$$
\begin{equation*}
Y_{n}=\sum_{\substack{I \subset\{1, \ldots, l\} \\|I|=n}}\left(\prod_{\substack{j \in I \\ k \in I^{c}}} \frac{\sigma\left(x_{j}-x_{k}+\mu ; \omega_{1}, \omega_{2}\right)}{\sigma\left(x_{j}-x_{k} ; \omega_{1}, \omega_{2}\right)}\right) \prod_{j \in I} \tau_{j}\left(\omega_{3}\right), \tag{1.1}
\end{equation*}
$$

where $\omega_{1}, \omega_{2} \in \mathbb{C} \backslash\{0\}$ are arbitrary such that $\omega_{1} / \omega_{2} \notin \mathbb{R}$, and $\omega_{3}, \mu \in \mathbb{C} \backslash\{0\}$ and the action of $\tau_{j}(\omega)$ is defined by $\left(\tau_{j}(\omega) f\right)\left(x_{1}, \ldots, x_{j}, \ldots, x_{n}\right)=f\left(x_{1}, \ldots, x_{j}-\omega, \ldots, x_{n}\right)$. It is shown that these operators are mutually commutative. The first result of this article is a generalization of the elliptic Ruijsenaars operators. We define

$$
\begin{equation*}
Y_{n}^{(p)}=\sum_{\substack{I \subset\{1, \ldots, l\} \\|I|=n}}\left(\prod_{\substack{j \in I \\ k \in I^{c}}} e^{\nu_{p}\left(x_{j}-x_{k}\right)} \frac{\sigma\left(x_{j}-x_{k}+\mu_{p} ; \omega_{q}, \omega_{r}\right)}{\sigma\left(x_{j}-x_{k} ; \omega_{q}, \omega_{r}\right)}\right) \prod_{j \in I} \tau_{j}\left(\omega_{p}\right), \tag{1.2}
\end{equation*}
$$

where $p, q, r \in \mathbb{Z} / 3 \mathbb{Z}$ are distinct, and for $p \in \mathbb{Z} / 3 \mathbb{Z}, \omega_{p} \in \mathbb{C} \backslash\{0\}$ are such that $\omega_{p} / \omega_{q} \notin \mathbb{R}$ if $p \neq q$. Put $\eta_{p q}=2 \zeta\left(\omega_{p} / 2 ; \omega_{p}, \omega_{q}\right)$ with Weierstrass' zeta function $\zeta$ and $a_{r}=\eta_{p q} \omega_{q}-\eta_{q p} \omega_{p}$, then $a_{r}= \pm 2 \pi i$. If $\nu_{p}, \mu_{p}$ satisfy three equations $\nu_{p} \omega_{q}+\mu_{p} \eta_{q r}=\nu_{q} \omega_{p}+\mu_{q} \eta_{p r}$ for distinct $p, q, r \in \mathbb{Z} / 3 \mathbb{Z}$, then all $Y_{n}^{(p)}$ are shown to be mutually commutative. For instance, these
equations are solved by

$$
\begin{gather*}
\nu_{1}=\frac{a_{1}\left(\nu_{3} \omega_{1}+\mu_{3} \eta_{12}\right)-\left(a_{2} \mu_{2}+a_{3} \mu_{3}\right) \eta_{32}}{a_{1} \omega_{3}},  \tag{1.3}\\
\nu_{2}=\frac{\nu_{3} \omega_{2}+\mu_{3} \eta_{21}-\mu_{2} \eta_{31}}{\omega_{3}},  \tag{1.4}\\
\mu_{1}=\frac{a_{2} \mu_{2}+a_{3} \mu_{3}}{a_{1}} \tag{1.5}
\end{gather*}
$$

where $\nu_{3}, \mu_{2}, \mu_{3}$ are regarded as free parameters. Although the discussion above is for $A$-type root system, the construction can be applied to arbitrary root systems.
The second result is a construction of a simultaneous eigenvector of the elliptic Ruijsenaars operators of type $B C_{1}$. We obtain an explicit meromorphic eigenvector described in terms of the elliptic hypergeometric integral. Note that for the $A$-type root system, some classes of eigenvectors are discussed in [1, 4, 6, 12, 15-18].

## 2. Affine Root Systems

We summarize some facts about affine root systems and affine Weyl groups $[2,3,7$, 8]. In this article, we will omit $A_{2 l}^{(2)}$-type root system because of simplicity, though it is straightforward. The notation and symbols are a little different from those in the previous papers $[9,10]$ in order to generalize the results.

Let $\Delta$ be the irreducible reduced finite root system of type $X_{l}$ in a complex vector space $V$ with $\operatorname{dim} V=l$ and the inner product $\langle\cdot, \cdot\rangle, I=\{1, \ldots, l\}$ a set of indices, $\Pi=\left\{\alpha_{i} \mid i \in I\right\} \subset V$ the set of simple roots, $\Pi^{\vee}=\left\{\alpha_{i}^{\vee} \mid i \in I\right\} \subset V$ the set of simple coroots, $Q$ and $Q^{\vee}$ the root and coroot lattices, $P$ and $P^{\vee}$ the weight and coweight lattices, $\left\{\Lambda_{i} \mid i \in I\right\}$ and $\left\{\Lambda_{i}^{\vee} \mid i \in I\right\}$ the fundamental weights and fundamental coweights such that $\left\langle\alpha_{i}, \Lambda_{j}^{\vee}\right\rangle=\left\langle\Lambda_{i}, \alpha_{j}^{\vee}\right\rangle=\delta_{i j}$. Then we have

$$
\begin{align*}
Q & =\bigoplus_{i \in I} \mathbb{Z} \alpha_{i} \subset P=\bigoplus_{i \in I} \mathbb{Z} \Lambda_{i} \subset V,  \tag{2.1}\\
Q^{\vee} & =\bigoplus_{i \in I} \mathbb{Z} \alpha_{i}^{\vee} \subset P^{\vee}=\bigoplus_{i \in I} \mathbb{Z} \Lambda_{i}^{\vee} \subset V . \tag{2.2}
\end{align*}
$$

The inner product $\langle\cdot, \cdot\rangle$ is normalized such that $\langle\alpha, \alpha\rangle=2$ for the longer roots $\alpha$. Let $\Delta_{+}$ and $\Delta_{-}$be the set of positive roots and negative roots respectively.

Let $\Delta_{s} \subset \Delta$ be the set of shorter roots and $\Delta_{l} \subset \Delta$ the set of longer roots. Let $r$ be the ratio of the square lengths of longer roots and shorter roots. Fix parameters $\gamma_{\alpha}$ for $\alpha \in \Delta$, such that in nontwisted case all $\gamma_{\alpha}=1$, and in twisted case, $\gamma_{\alpha}=r$ if $\alpha \in \Delta_{l}, \gamma_{\alpha}=1$ otherwise. Let $\widehat{V}=V \oplus \mathbb{C} \delta$ with $\left\langle\alpha_{i}, \delta\right\rangle=\langle\delta, \delta\rangle=0$ and its linear extension. Then the associated affine root system $\widehat{\Delta} \subset \widehat{V}$ is written as

$$
\begin{equation*}
\widehat{\Delta}=\left\{\alpha+n \gamma_{\alpha} \delta \mid \alpha \in \Delta, n \in \mathbb{Z}\right\} \tag{2.3}
\end{equation*}
$$

Let $\widehat{\Delta}_{+}$and $\widehat{\Delta}_{-}$be the set of positive affine roots and negative affine roots respectively. We denote by $\bar{v}$ for $v \in \widehat{V}$ the natural projection on $V$.

For $\alpha \in \widehat{\Delta}$, let $s_{\alpha}$ be a reflection defined by

$$
\begin{equation*}
s_{\alpha}(v):=v-\langle\alpha, v\rangle \alpha^{\vee}, \quad v \in V \tag{2.4}
\end{equation*}
$$

The Weyl group $W$ is generated by the fundamental reflections $\left\{s_{i}:=s_{\alpha_{i}} \mid i \in I\right\}$ on $V$ and the affine Weyl group $\widehat{W}$ is generated by $\left\{s_{i} \mid i \in \widehat{I}\right\}$, where $\widehat{I}=I \cup\{0\}$ and $\alpha_{0}=\delta-\theta$ with $\theta$ the highest root in nontwisted case and the highest short root in twisted case.
The defining relations are given by $s_{i}^{2}=i d$ and the Coxeter relations:

$$
\begin{equation*}
\left(s_{i} s_{j}\right)^{m_{i j}}=i d, \quad \text { for } i \neq j \in \widehat{I} \tag{2.5}
\end{equation*}
$$

where $m_{i j}=2$ if $\alpha_{i}$ and $\alpha_{j}$ are disconnected in the Dynkin diagram and $m_{i j}=3,4,6$ if $1,2,3$ lines respectively connect $\alpha_{i}$ and $\alpha_{j}$. For $\mu \in V$, we define endomorphisms $\tau_{\mu}$ of the vector space $V$ by

$$
\begin{equation*}
\tau_{\mu}(\lambda):=\lambda-\langle\lambda, \mu\rangle \delta . \tag{2.6}
\end{equation*}
$$

Let $M:=\mathbb{Z}\left(W \cdot \theta^{\vee}\right) \subset V$. For an arbitrary lattice $L$, we denote by $T_{L}$ the corresponding group of translations of $L$. Then one sees that $\widehat{W}$ is the semidirect product $\widehat{W}=W \ltimes T_{M}$. Let $\widetilde{M}:=\left\{\lambda \in V \mid \alpha \in \Delta,\langle\alpha, \lambda\rangle \in \gamma_{\alpha} \mathbb{Z}\right\}$. The extended affine Weyl group $\widetilde{W}$ is defined by the semidirect product $\widetilde{W}:=W \ltimes T_{\widetilde{M}}$. Let $\Omega$ be the subgroup of $\widetilde{W}$ which stabilizes the affine Weyl chamber $C$. Then one sees that $\widetilde{W}$ is isomorphic to the semidirect product $\widehat{W} \rtimes \Omega$. Here are the explicit description of $\widetilde{M}$ and its canonical basis $\left\{\lambda_{i} \mid i \in I\right\}$ :

$$
\widetilde{M}=\left\{\begin{array}{ll}
P^{\vee}, & \text { nontwisted case, }  \tag{2.7}\\
P, & \text { twisted case },
\end{array} \quad \lambda_{i}= \begin{cases}\Lambda_{i}^{\vee}, & \text { nontwisted case }, \\
\Lambda_{i}, & \text { twisted case }\end{cases}\right.
$$

We also use $\widetilde{M}_{-}:=\oplus_{i \in I} \mathbb{Z}_{\leq 0} \lambda_{i}$.
The length $\ell(w)$ of $w \in \widetilde{W}$ is defined by the length $\ell$ of a reduced decomposition:

$$
\begin{equation*}
w=s_{i_{1}} \ldots s_{i_{\ell}} \omega, \quad i_{k} \in \widehat{I}, \omega \in \Omega \tag{2.8}
\end{equation*}
$$

It is equivalent to the number of the negative roots made positive by $\hat{w}$ :

$$
\begin{equation*}
\ell(\hat{w}):=\left|\Delta_{\hat{w}}\right|, \quad \Delta_{\hat{w}}:=\widehat{\Delta}_{+} \cap \hat{w} \widehat{\Delta}_{-} \tag{2.9}
\end{equation*}
$$

The set $\Delta_{\hat{w}}$ is explicitly described as $\Delta_{\hat{w}}=\left\{\alpha^{(1)}=\alpha_{i_{1}}, \alpha^{(2)}=s_{i_{1}}\left(\alpha_{i_{2}}\right), \ldots, \alpha^{(\ell)}=\right.$ $\left.w s_{i_{\ell}}\left(\alpha_{i_{\ell}}\right)\right\}$. By definition, $\Delta_{\hat{w}}$ is independent of reduced expressions. One sees that $\Omega=\{\omega \in \widehat{W} \mid \ell(\omega)=0\}$. A weight $\lambda \in \widetilde{M}$ is said to be minuscule if $\Delta_{\tau_{-\lambda}} \subset \Delta_{+}$.

In the following, we use the constants

$$
\begin{equation*}
\rho_{x}:=\sum_{i \in I} x_{\alpha_{i}} \Lambda_{i}=\frac{1}{2} \sum_{\alpha \in \Delta_{+}} x_{\alpha} \alpha, \tag{2.10}
\end{equation*}
$$

where $x_{\alpha} \in \mathbb{C}$ which depends only on the length of roots,
We shall define the root algebras after Cherednik [3,10].
Definition 2.1. Root algebra $\mathcal{R}$ is generated by independent variables $\left\{R_{\alpha} \mid \alpha \in \widehat{\Delta}\right\}$ and $\left\{\tau_{\lambda} \mid \lambda \in \widetilde{M}\right\}$ with the following defining relations:

$$
\begin{gather*}
\underbrace{R_{w \alpha_{i}} R_{w s_{i} \alpha_{j}} R_{w s_{i} s_{j} \alpha_{i}} \cdots}_{m_{i j} \text { factors }}=\underbrace{R_{w \alpha_{j}} R_{w s_{j} \alpha_{i}} R_{w s_{j} s_{i} \alpha_{j}} \cdots}_{m_{i j} \text { factors }}, \quad \text { for } w \in \widetilde{W}  \tag{2.11}\\
\tau_{\lambda} R_{\alpha}=R_{\tau_{\lambda} \alpha} \tau_{\lambda},  \tag{2.12}\\
\tau_{\lambda} \tau_{\lambda^{\prime}}=\tau_{\lambda+\lambda^{\prime}} . \tag{2.13}
\end{gather*}
$$

Theorem 2.2 (Cherednik).

1. There exists a unique set $\left\{R_{w} \mid w \in \widetilde{W}\right\} \subset \mathcal{R}$ satisfying the relations:

$$
\begin{equation*}
R_{v w}=R_{v}{ }^{v} R_{w}, \quad R_{s_{i}}=R_{\alpha_{i}} \quad(i \in \widehat{I}), \quad R_{\omega}=1 \tag{2.14}
\end{equation*}
$$

where $\omega \in \Omega, v, w \in \widetilde{W}$ and $\ell(v w)=\ell(v)+\ell(w)$, and ${ }^{v}\left(R_{\alpha_{1}} \ldots R_{\alpha_{i}}\right)=R_{v \alpha_{1}} \ldots R_{v \alpha_{i}}$
2. We have the $R$-matrix for $w \in \widetilde{W}$ and its arbitrary reduced decomposition $w=$ $s_{i_{1}} \ldots s_{i_{\ell}} \omega$ as

$$
\begin{align*}
& R_{w}=R_{\alpha(1)} \ldots R_{\alpha^{(\ell)}}, \\
& \alpha^{(1)}=\alpha_{i_{1}}, \quad \alpha^{(2)}=s_{i_{1}}\left(\alpha_{i_{2}}\right), \quad \ldots, \quad \alpha^{(\ell)}=w s_{i_{\ell}}\left(\alpha_{i_{\ell}}\right) \in \Delta_{w} . \tag{2.15}
\end{align*}
$$

Theorem 2.3. The subalgebra $\mathcal{S} \subset \mathcal{R}$ generated by $\left\{Y^{\lambda}:=R_{\tau_{\lambda}} \tau_{\lambda} \mid \lambda \in \widetilde{M}_{-}\right\}$forms a commutative algebra and is generated by $\left\{Y^{-\lambda_{i}} \mid i \in I\right\}$.

## 3. Representation and Difference Operators

Let ${\underset{\alpha}{\alpha}}_{(1)}^{(1)}=\gamma_{\alpha}$ and $\gamma_{\alpha}^{(2)}, \gamma_{\alpha}^{(3)}$ be taken similarly as $\gamma_{\alpha}^{(1)}$. Accordingly let $\widetilde{M}^{(1)}=\widetilde{M}$ and $\widetilde{M}^{(2)}, \widetilde{M}^{(3)}$ be taken similarly. Here $\gamma_{\alpha}^{(i)}$ and $\gamma_{\alpha}^{(j)}$ may differ if $i \neq j$. Let $\mathcal{M}$ be the set of meromorphic functions on $V$. To define the action of $\widetilde{W}$ on $\mathcal{M}$, it is sufficient to specify the action of $s_{i}$ for $i \in I$ and $\tau_{\lambda}$ for $\lambda \in \widetilde{M}^{(1)}$. For $f \in \mathcal{M}$, we define

$$
\begin{equation*}
s_{i}(f)(v)=f\left(s_{i}(v)\right), \quad \tau_{\lambda}(f)(v)=\tau_{\lambda}^{(1)}(f)(v)=f\left(v-\omega_{1} \lambda\right) \tag{3.1}
\end{equation*}
$$

Fix $\xi^{(1)}, \zeta^{(1)} \in V$ and let $\mu_{\alpha}^{(1)}, \nu_{\alpha}^{(1)}$ be constants depending only on the length of roots. Then one can check that (3.1) and

$$
\begin{equation*}
\left(R_{\alpha} f\right)(v):=H_{\alpha}\left(\mu_{\alpha}^{(1)}, \nu_{\alpha}^{(1)}\right) f(v)-H_{\alpha}\left(\left\langle\xi^{(1)}, \alpha^{\vee}\right\rangle,\left\langle\zeta^{(1)}, \alpha^{\vee}\right\rangle\right) f\left(s_{\alpha} v\right), \tag{3.2}
\end{equation*}
$$

satisfy the defining relations of the root algebra [11], where $H_{\alpha}(\eta, \kappa)$ is a meromorphic function defined by

$$
\begin{equation*}
H_{\alpha}(\eta, \kappa)(v):=e^{\kappa \alpha(v)} \frac{\sigma\left(\mu_{\alpha}^{(1)} ; \gamma_{\alpha}^{(2)} \omega_{2}, \gamma_{\alpha}^{(3)} \omega_{3}\right)}{\sigma\left(\eta ; \gamma_{\alpha}^{(2)} \omega_{2}, \gamma_{\alpha}^{(3)} \omega_{3}\right)} \frac{\sigma\left(\alpha(v)+\eta ; \gamma_{\alpha}^{(2)} \omega_{2}, \gamma_{\alpha}^{(3)} \omega_{3}\right)}{\sigma\left(\alpha(v) ; \gamma_{\alpha}^{(2)} \omega_{2}, \gamma_{\alpha}^{(3)} \omega_{3}\right)}, \tag{3.3}
\end{equation*}
$$

and a root $\alpha$ acts on $V$ as an affine linear functional $\alpha(v)=\langle\alpha, v\rangle+n \omega_{1}$ for $v \in V$ and $\alpha=\alpha^{\prime}+n \delta, \alpha^{\prime} \in \Delta$. Then we have the following theorems $[9,10]$

Theorem 3.1. Let $\mathcal{V}:=\mathcal{M}^{W}$, the $W$-invariant subspace of $\mathcal{M}$ and let $\xi^{(1)}=-\rho_{\mu^{(1)}}$, $\zeta^{(1)}=-\rho_{\nu^{(1)}}$. Then $Y_{\lambda}^{(1)}:=Y^{\lambda} \in \operatorname{End}_{\mathbb{C}} \mathcal{V}$.

Theorem 3.2. Let $(-\lambda) \in \widetilde{M}^{(1)}$ be minuscule. Then we have

$$
\begin{equation*}
\left.Y^{\lambda}\right|_{\mathcal{V}}=\left.\frac{1}{\left|W_{\lambda}\right|} \sum_{w \in W} w\left(\prod_{\substack{\alpha \in \Delta_{+} \\\langle\lambda, \alpha\rangle=-\gamma_{\alpha}}} H_{\alpha}\left(\mu_{\alpha}^{(1)}, \nu_{\alpha}^{(1)}\right) \tau_{\lambda}^{(1)}\right)\right|_{\mathcal{V}} \tag{3.4}
\end{equation*}
$$

where $W_{\lambda}$ is the stabilizer of $\lambda$ in $W$.
Let $Y_{\lambda}^{(2)}$ be the operator obtained by changing the role of the indices 1 and 2 in the construction of $Y_{\lambda}^{(1)}$ and $Y_{\lambda}^{(3)}$ be obtained in the same manner. In the following, we assume that $\mu_{\alpha}^{(j)} \eta_{k l}+\nu_{\alpha}^{(j)} \omega_{k}=\mu_{\alpha}^{(k)} \eta_{j l}+\nu_{\alpha}^{(k)} \omega_{j}$ where $j \neq k \neq l \neq j$. Then one sees that $Y_{\lambda}^{(j)}$ for $\lambda \in \widetilde{M}^{(j)}$ are commutative by Theorem 2.3 for a fixed $j$. Let $\xi^{(j)}=-\rho_{\mu^{(j)}}$ and $\zeta^{(j)}=-\rho_{\nu^{(j)}}$ for $j=1,2,3$.

Theorem 3.3. Let $j, k \in\{1,2,3\}$ and $-\lambda \in \widetilde{M}^{(j)},-\nu \in \widetilde{M}^{(k)}$. Then $Y_{\lambda}^{(j)}$ and $Y_{\nu}^{(k)}$ are commutative.

For minuscule weights $-\lambda$, the periodicity of the coefficients is easily obtained since the explicit forms of the operators $Y^{\lambda}$ are calculated [10]. Hence the commutativity of the operators for $A$-type root system follows because all the fundamental weights are minuscule in this root system. However $Y^{\lambda}$ for general $\lambda$ is complicated and the proof of the commutativity requires a further investigation.

## 4. $B C_{1}$-type Operators and Eigenvector

Generally, it is very difficult to construct explicit eigenvectors of the elliptic Ruijsenaars operators. However we can construct simultaneous eigenvectors in the $B C_{1}$ root system. In this root system, there are three mutually commutative operators.

Let $\omega_{1}, \omega_{2}, \mu_{0}, \ldots, \mu_{6} \in \mathbb{C}$ such that $\Im \omega_{1}, \omega_{2}>0$ and $2 \mu_{0}+\mu_{1}+\cdots+\mu_{6}+\omega_{1}+\omega_{2}=0$. Then the operators are given by

$$
\begin{align*}
& \quad \vartheta\left(x-\mu_{0}+\omega_{1} ; \omega_{2}\right) \prod_{j=0}^{6} \vartheta\left(x-\mu_{j} ; \omega_{2}\right)  \tag{4.1}\\
& Y^{(1)}=e^{2 \pi i x} \frac{\vartheta\left(2 x ; \omega_{2}\right) \vartheta\left(2 x+\omega_{1} ; \omega_{2}\right)}{}\left(\tau\left(\omega_{1}\right)-1\right), \\
&+(x \leftrightarrow-x) \\
& Y^{(2)}=e^{2 \pi i x} \frac{\vartheta\left(x-\mu_{0}+\omega_{2} ; \omega_{1}\right) \prod_{j=0}^{6} \vartheta\left(x-\mu_{j} ; \omega_{1}\right)}{\vartheta\left(2 x ; \omega_{1}\right) \vartheta\left(2 x+\omega_{2} ; \omega_{1}\right)}\left(\tau\left(\omega_{2}\right)-1\right),  \tag{4.2}\\
&+(x \leftrightarrow-x) \\
& Y^{(3)}=\tau(1)+\tau(-1)-2, \tag{4.3}
\end{align*}
$$

where $\vartheta$ denotes the Jacobi odd theta function. For these commutative difference operators we can construct an explicit simultaneous eigenvector:

Theorem 4.1. Let

$$
\begin{align*}
\phi(c) & =\frac{\prod_{i=1}^{7} \Gamma\left(p q c_{0} / c_{i} ; p, q\right)}{\Gamma\left(p q c_{0} ; p, q\right) \prod_{1 \leq i<j \leq 7} \Gamma\left(p q c_{0} / c_{i} c_{j} ; p, q\right)}  \tag{4.4}\\
& \times \int_{C} \frac{\Gamma\left(p q /\left(p q c_{0}\right)^{1 / 2} z^{ \pm 1} ; p, q\right) \prod_{j=1}^{7} \Gamma\left(\left(p q c_{0}\right)^{1 / 2} / c_{j} z^{ \pm 1} ; p, q\right)}{\Gamma\left(z^{2}, z^{-2} ; p, q\right)} \frac{d z}{z},
\end{align*}
$$

with

$$
\begin{gathered}
c_{0}=e^{2 \pi i\left(\mu_{0}-\mu_{1}-\omega_{1}-\omega_{2}\right)}, \quad c_{1}=e^{2 \pi i\left(\mu_{0}+\mu_{2}\right)}, \quad c_{2}=e^{2 \pi i\left(\mu_{0}+\mu_{3}\right)}, \\
c_{3}=e^{2 \pi i\left(\mu_{0}+\mu_{4}\right)}, \quad c_{4}=e^{2 \pi i\left(\mu_{0}+\mu_{5}\right)}, \quad c_{5}=e^{2 \pi i\left(\mu_{0}+\mu_{6}\right)}, \\
c_{6}=e^{2 \pi i\left(-x-\mu_{1}\right)}, \quad c_{7}=e^{2 \pi i\left(x-\mu_{1}\right)}, \\
p=e^{2 \pi i \omega_{1}}, \quad q=e^{2 \pi i \omega_{2}},
\end{gathered}
$$

and $C$ a closed circle taken appropriately, and $\Gamma(z ; p, q)$ elliptic Gamma function defined by

$$
\Gamma(z ; p, q)=\frac{\left(p q z^{-1} ; p, q\right)_{\infty}}{(z ; p, q)_{\infty}}
$$

Then $\phi(c)$ is a simultaneous eigenvector of $Y^{(j)}$ with the eigenvalues

$$
\begin{aligned}
& E^{(1)}=e^{-2 \pi i \mu_{0}} \prod_{i=1}^{6} \vartheta\left(\mu_{0}+\mu_{i} ; \omega_{2}\right), \quad E^{(3)}=0, \\
& E^{(2)}=e^{-2 \pi i \mu_{0}} \prod_{i=1}^{6} \vartheta\left(\mu_{0}+\mu_{i} ; \omega_{1}\right) .
\end{aligned}
$$

An integral analogue of Bailey transformation and the other properties of the integral factor of (4.4) are investigated in [13].

## 5. Special Cases

In this section, we will clarify the relation between $\phi(c)$ and the elliptic hypergeometric series ${ }_{12} V_{11}$ introduced by Frenkel and Turaev [5].

Theorem 5.1. If $c_{j}=p^{-M} q^{-N}$ for some $j$, then $\phi(c)$ (4.4) splits into two elliptic hypergeometric series,

$$
\begin{equation*}
\phi(c)=\tilde{\phi}\left(c_{0}, \ldots, q^{-N}, \ldots, c_{7} ; p\right) \times \tilde{\phi}\left(c_{0}, \ldots, p^{-M}, \ldots, c_{7} ; q\right), \tag{5.1}
\end{equation*}
$$

where

$$
\begin{equation*}
\tilde{\phi}\left(c_{0}, \ldots, c_{7} ; r\right)=\sum_{k=0}^{\infty} \frac{\left\langle(p q)^{2 k} c_{0} ; r\right\rangle}{\left\langle c_{0} ; r\right\rangle} \frac{\left\langle\left\langle c_{0} ; r\right\rangle\right\rangle_{k}}{\langle\langle p q ; r\rangle\rangle_{k}} \prod_{i=1}^{7} \frac{\left\langle\left\langle c_{i} ; r\right\rangle_{k}\right.}{\left\langle\left\langle p q c_{0} / c_{i} ; r\right\rangle\right\rangle_{k}}, \tag{5.2}
\end{equation*}
$$

and $\langle\langle u ; r\rangle\rangle_{k}=\langle u ; r\rangle \cdots\left\langle u(p q)^{k-1} ; r\right\rangle$ with $\langle u, r\rangle=u^{-1 / 2}(u ; r)_{\infty}\left(u^{-1} r ; r\right)_{\infty}$.
Now we state the relation between $\phi(c)$ and ${ }_{12} V_{11}$ which is defined by

$$
{ }_{12} V_{11}\left(b_{0} ; b_{1}, \ldots, b_{7} ; r\right)=\sum_{k=0}^{\infty} \frac{\left\langle q^{2 k} b_{0} ; r\right\rangle}{\left\langle b_{0} ; r\right\rangle} \frac{\left\langle b_{0} ; r\right\rangle_{k}}{\langle r ; r\rangle_{k}} \prod_{i=1}^{7} \frac{\left\langle b_{i} ; r\right\rangle_{k}}{\left\langle r b_{0} / b_{i} ; r\right\rangle_{k}},
$$

where one of $b_{i}$ should be $r^{-N}$ for the series to terminate and $\langle u ; r\rangle_{k}=\langle u ; r\rangle \cdots\left\langle u r^{k-1} ; r\right\rangle$. The following is shown by a direct calculation.

## Theorem 5.2.

$$
\tilde{\phi}\left(c_{0}, \ldots, c_{7} ; p\right)={ }_{12} V_{11}\left(p c_{0} ; c_{1}, \ldots, p c_{j}, \ldots, c_{7} ; p\right),
$$

which is independent of the choice of $j$.

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